



BROWN UNIVERSITY

PROVIDENCE, RHODE ISLAND

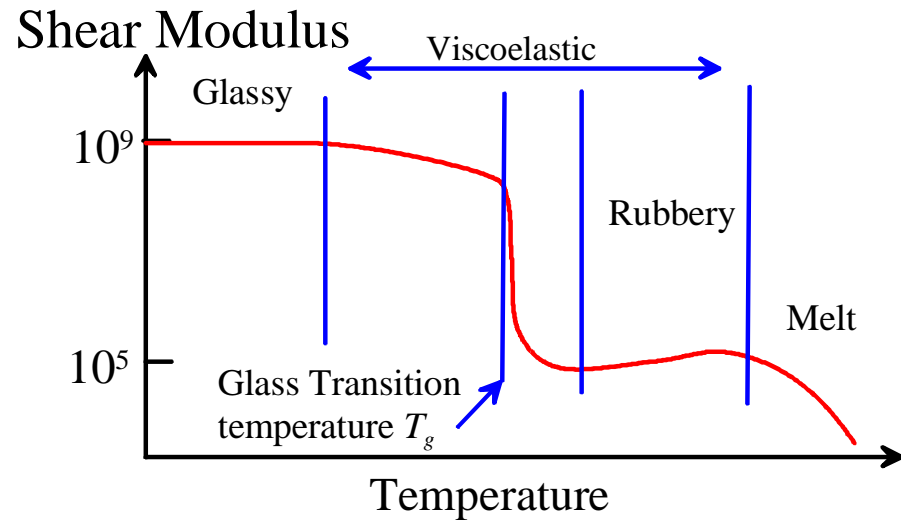
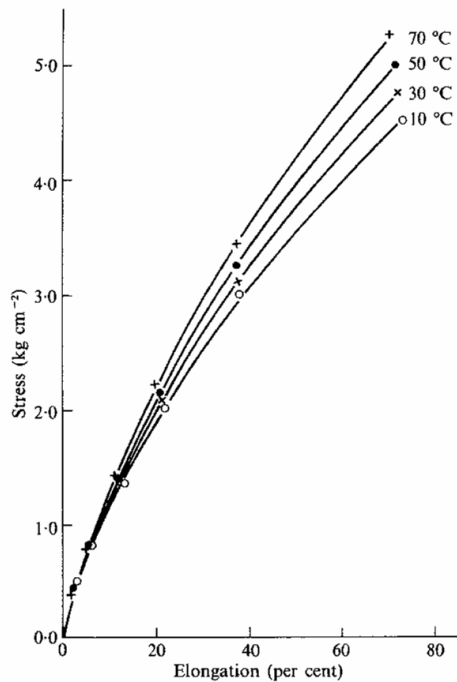


Division of Engineering  
Brown University

# Atomic scale mechanisms of stress production in elastomers

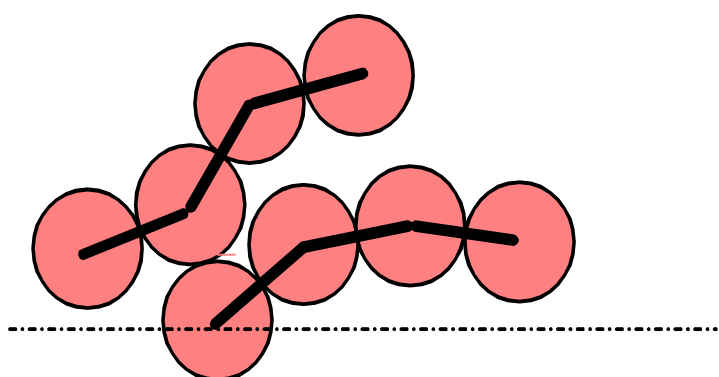
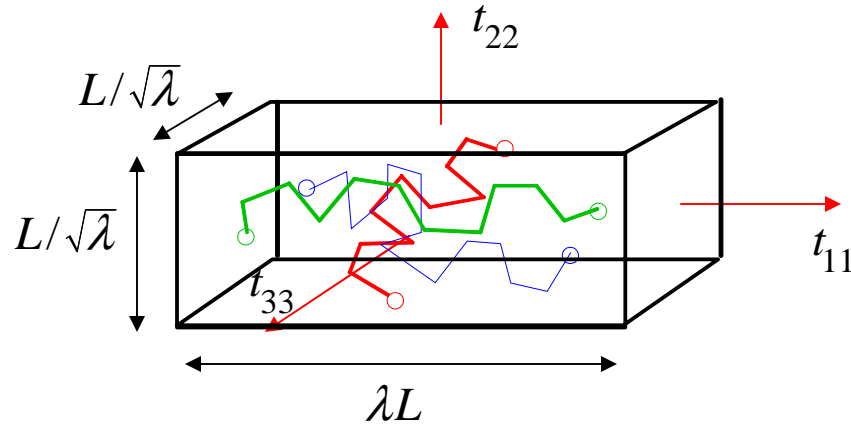
A.F. Bower, J.H. Weiner  
NSF Nanomechanics Workshop  
July 2004

# Constitutive response of elastomers



- Rubbers strongly resist volume changes. Bulk modulus comparable to metals or covalently bonded solids;
- Rubbers are very compliant in shear – shear modulus is several orders of magnitude lower than metals or covalently bonded solids
- Shear response is strongly temperature dependent: the material becomes stiffer as it is heated, in sharp contrast to metals;
- When stretched, rubber gives off heat.

# Classical theory of rubber elasticity



Mean stress:  $p = (t_{11} + t_{22} + t_{33})/3$  due to repulsive excluded volume interactions (fluid-like response)

Deviatoric stress:  $\sigma_{11} = t_{11} - p$  due to entropy of chain network

## Affine Network Model (freely jointed chains without EV)

Helmholtz free energy  $dW = dU - TdS = \sigma_{11} \frac{d\lambda}{\lambda}$

Entropy of Gaussian chain network  $S = S_0 - k \frac{N_c}{2\nu} (\lambda^2 + 2/\lambda)$

Deviatoric stress:  $\sigma_{11} = kT \frac{N_c}{\nu} (\lambda^2 - 1/\lambda)$

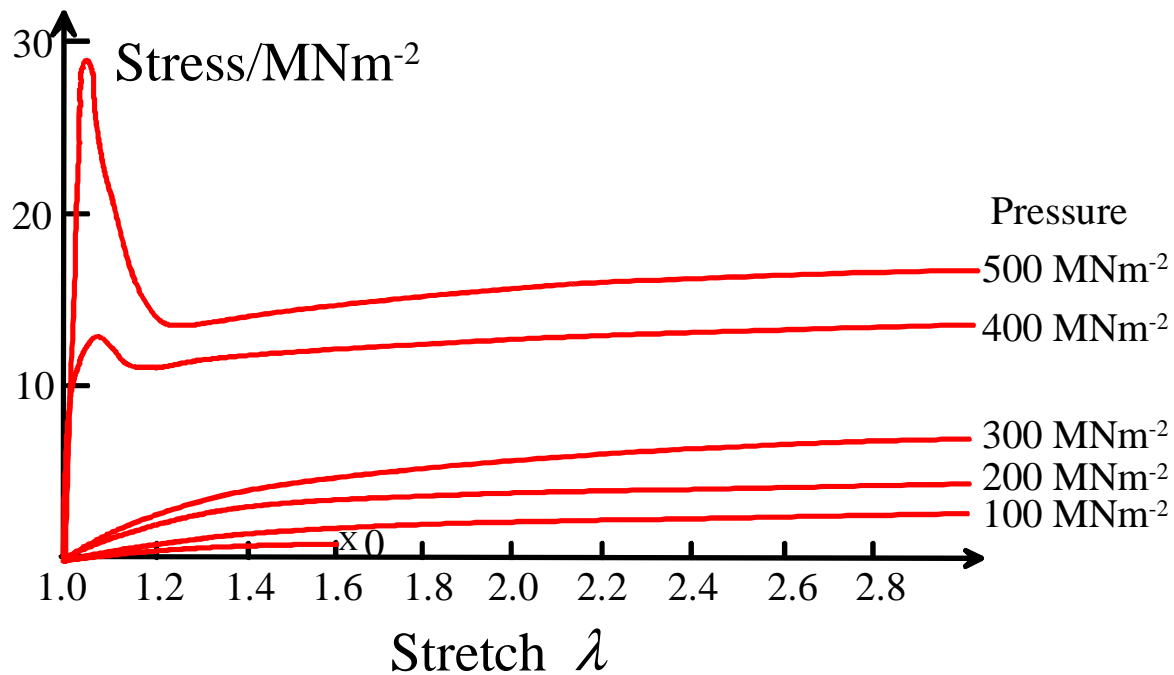
Chain force:  $f_i = \frac{3kT}{N_b a^2} r_i$



# The influence of pressure on deviatoric stress

D.L. Quested, K.D. Pae, J.L. Sheinbein and B.A. Newman,  
J. Appl. Phys, **52**, (10), 5977 (1981)

## Uniaxial deformation of solithane under confining pressure



Pressure dependence cannot be explained using classical theory of rubber elasticity



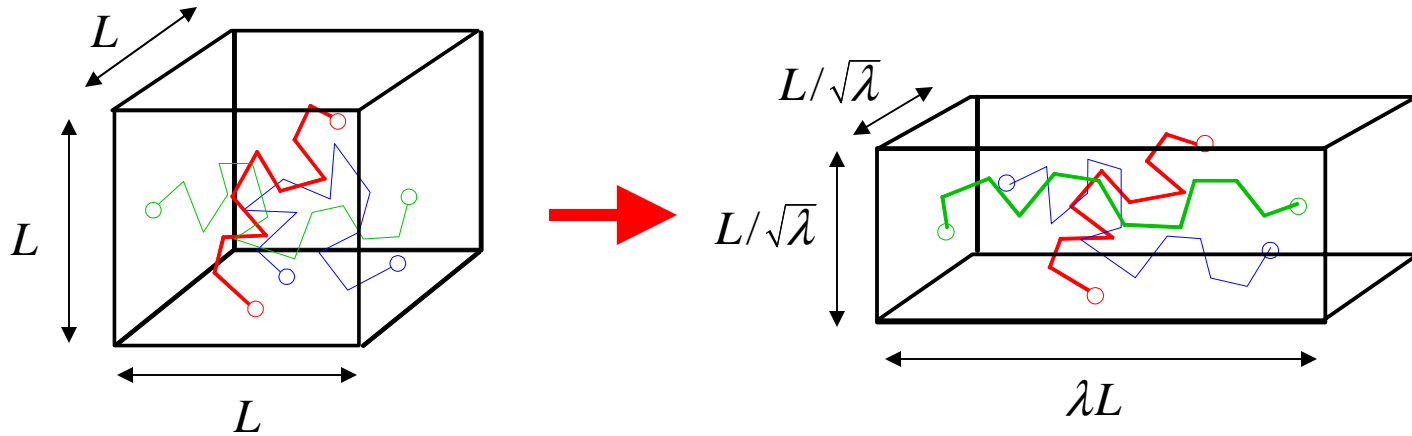
## Outline

- Re-examine atomic scale deformation of elastomers using a model that includes Excluded Volume (EV) interactions
- Deviatoric stress is caused by deformation induced anisotropy of chain network
- Repulsive EV interactions give the dominant contribution to both deviatoric and mean stress
- Pressure and deviatoric stress are intimately connected at atomic scale
- Pressure-deviatoric stress relation

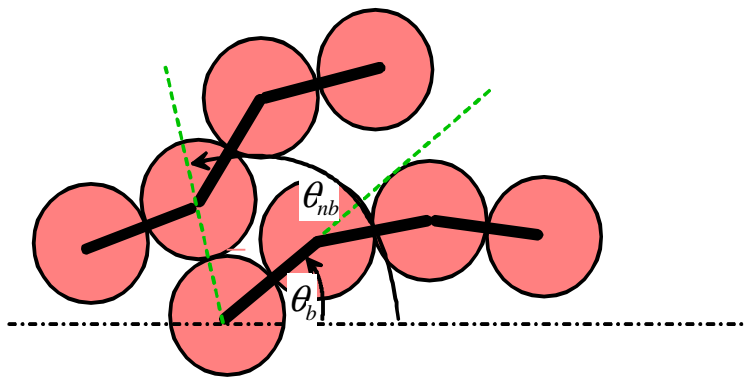
$$\sigma \approx \sigma_b + \sigma_{nbr} = \left( D\Pi_b - AB\Pi_{nbr} \right) \frac{F}{N_b} (\lambda^2 - \lambda^{-1})$$



# Molecular model of atomic-scale stress generation



Reduced density  $\rho^* = N\sigma_{LJ}^3 / L^3$



Covalent bonds

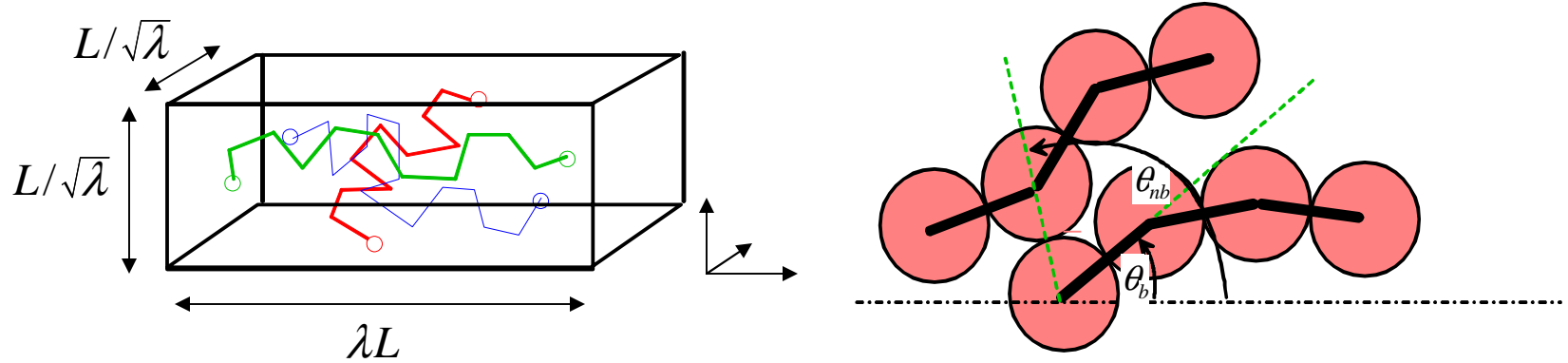
$$u_b(r) = \frac{1}{2} \kappa (r - a)^2$$

Noncovalent interactions

$$u_{nb}(r) = 4\epsilon_{LJ} \left[ \left( \frac{\sigma_{LJ}}{r} \right)^{12} - \left( \frac{\sigma_{LJ}}{r} \right)^6 \right]$$



# Stress computations



**Virial Theorem**

$$vt_{ij} = -NkT\delta_{ij} + \sum_{\alpha} \langle f_i^{\alpha} r_j^{\alpha} \rangle$$

**Normalized difference stress**

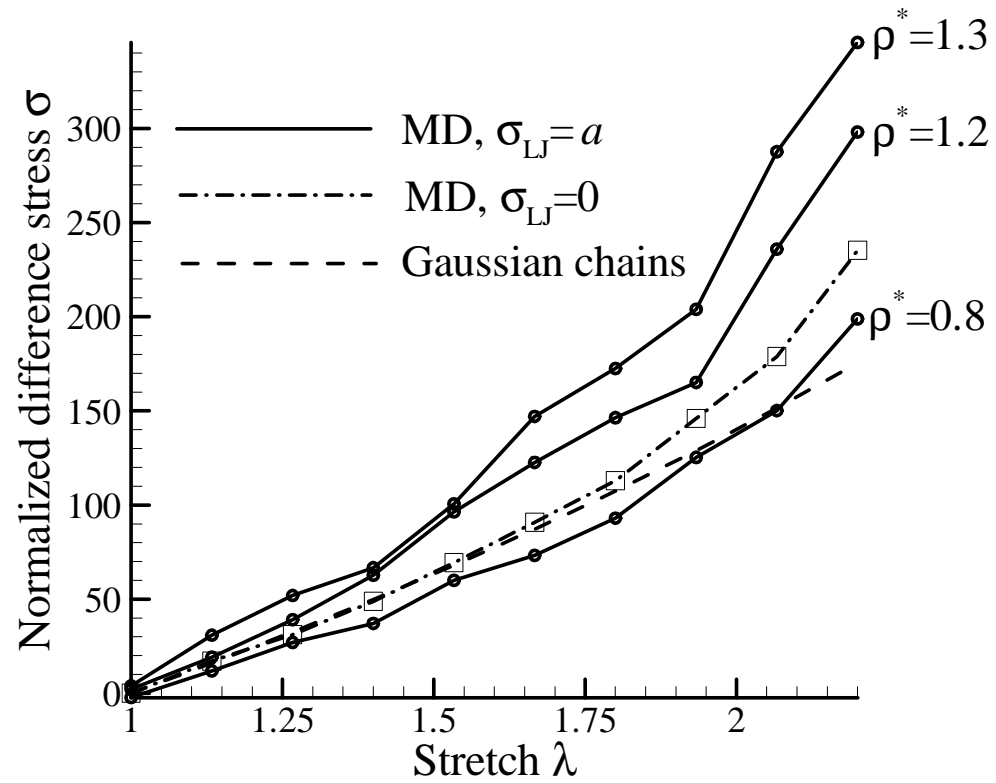
$$\begin{aligned} \sigma &= \frac{v}{kT} \{t_{11} - (t_{22} + t_{33})/2\} \\ &= \frac{1}{kT} \sum_{\alpha} \langle f_1^{\alpha} r_1^{\alpha} - (f_2^{\alpha} r_2^{\alpha} + f_3^{\alpha} r_3^{\alpha})/2 \rangle \end{aligned}$$

**Normalized mean stress**

$$\Pi = \frac{v}{kT} \{t_{11} + t_{22} + t_{33}\}/3 = -N + \frac{1}{3kT} \sum_{\alpha} \langle f_i^{\alpha} r_i^{\alpha} \rangle$$

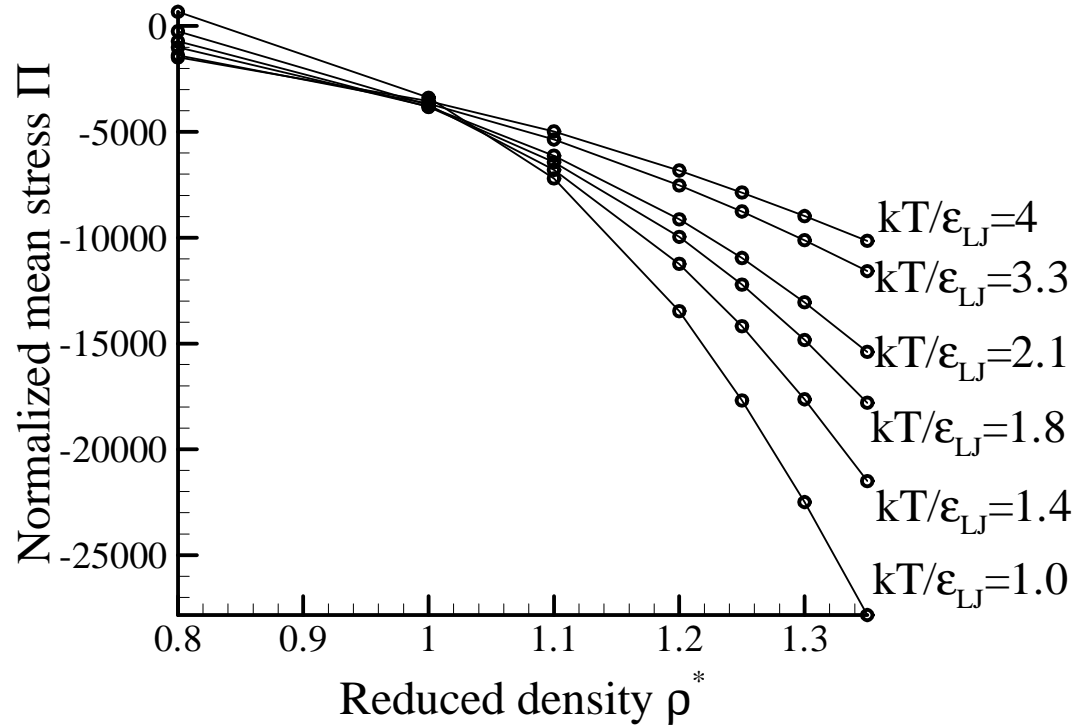


## Difference Stress-stretch relations





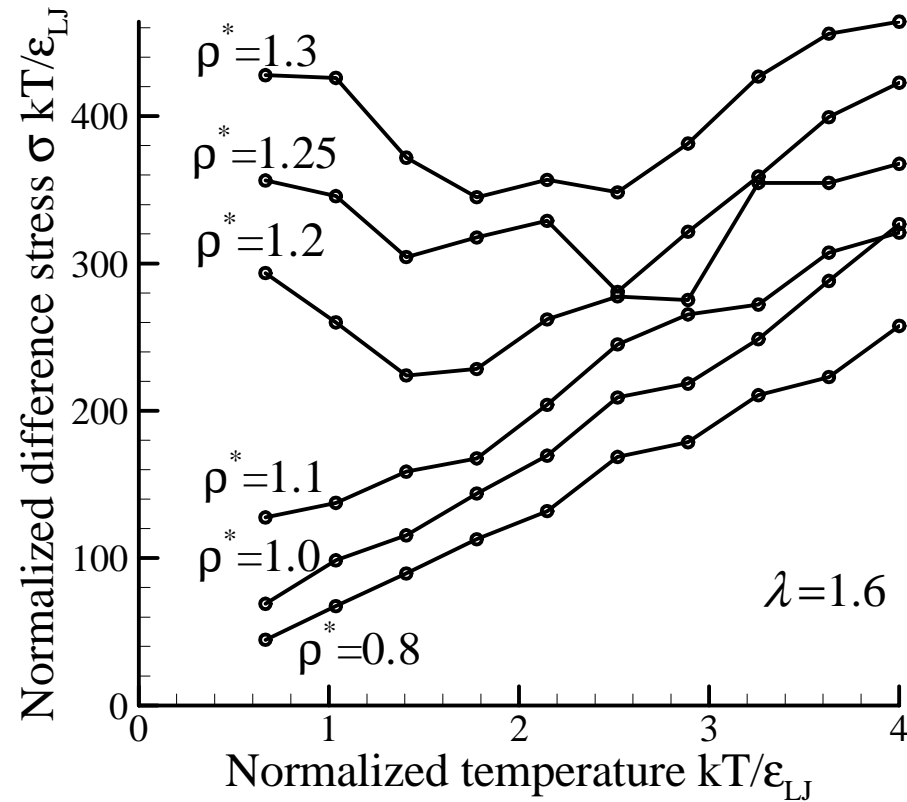
## Mean Stress – reduced density relations





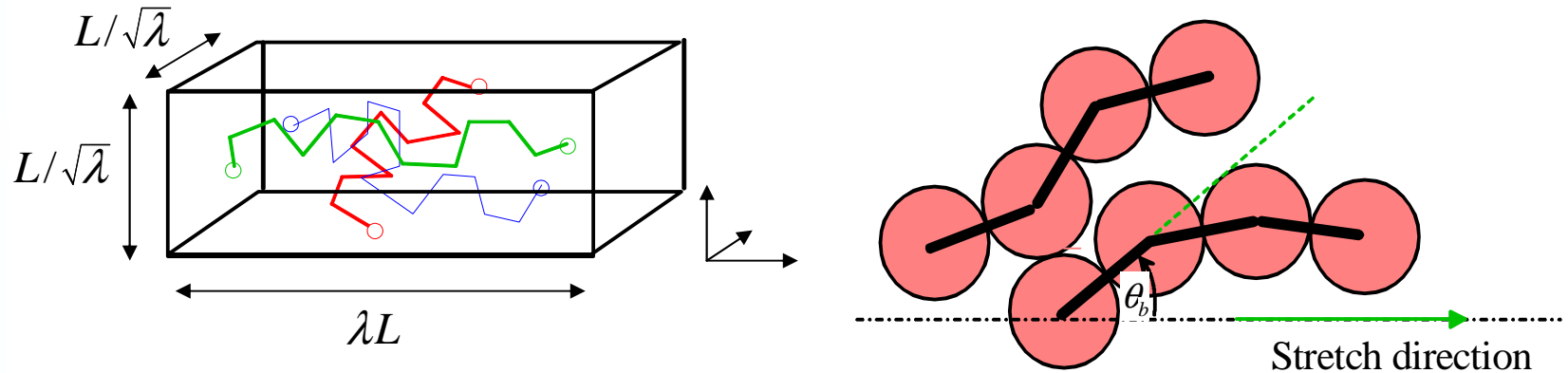
# Temperature dependence of stress

$$\frac{\sigma kT}{\epsilon_{LJ}} = \frac{v}{\epsilon_{LJ}} \{t_{11} - (t_{22} + t_{33})/2\}$$





# Relation between stress and bond anisotropy



## Bond orientation anisotropy measure

$$\langle P_2(\theta_b) \rangle = \frac{1}{N_c N_b} \sum_{\alpha \in b} \langle (3 \cos^2 \theta_\alpha - 1) / 2 \rangle$$

$\langle P_2(\theta_b) \rangle = 0$  for undeformed (isotropic) system.

$\langle P_2(\theta_b) \rangle > 0$  (tension)

$\langle P_2(\theta_b) \rangle < 0$  (compression)

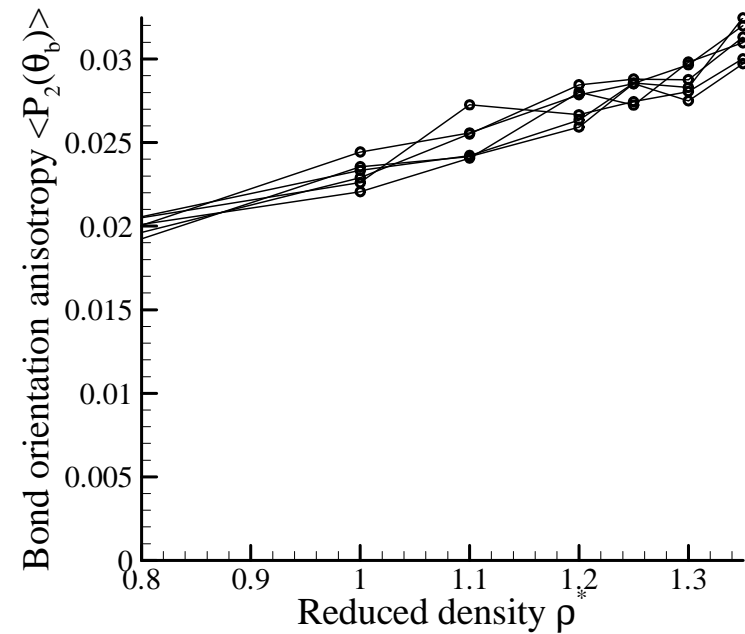
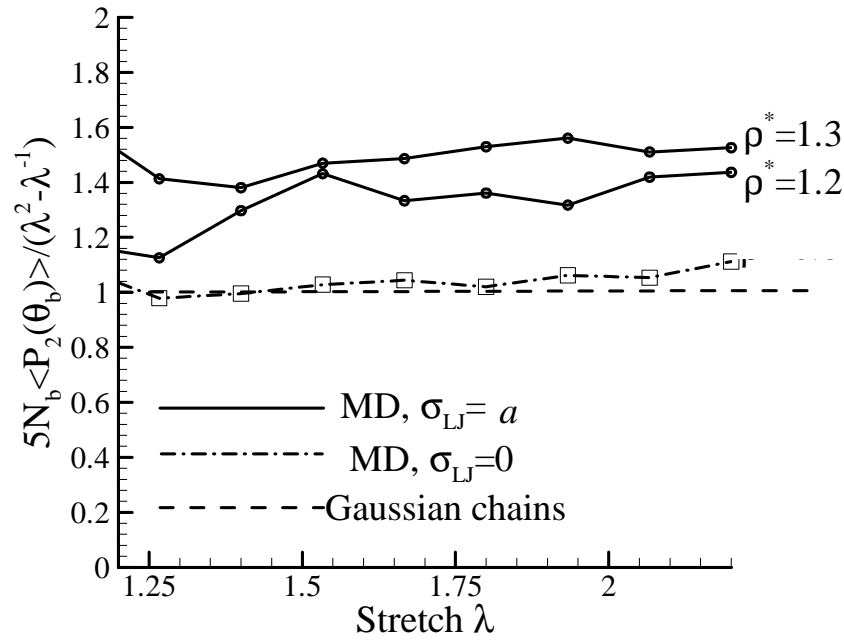
### For ideal chain network

$$\langle P_2(\theta_b) \rangle = \frac{1}{5N_b} (\lambda^2 - \lambda^{-1}) \quad \text{Difference stress per atom. } \frac{\sigma}{N} = \frac{1}{(N_b + 1)} (\lambda^2 - \lambda^{-1})$$

$$\frac{\sigma}{N \langle P_2(\theta_b) \rangle} = \frac{5N_b}{(N_b + 1)} = C$$



## Relation between stress and bond anisotropy



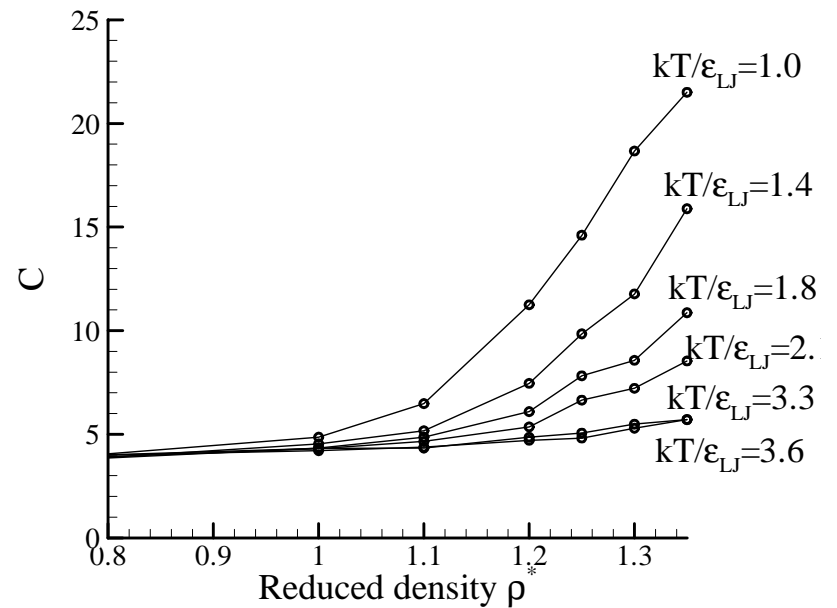
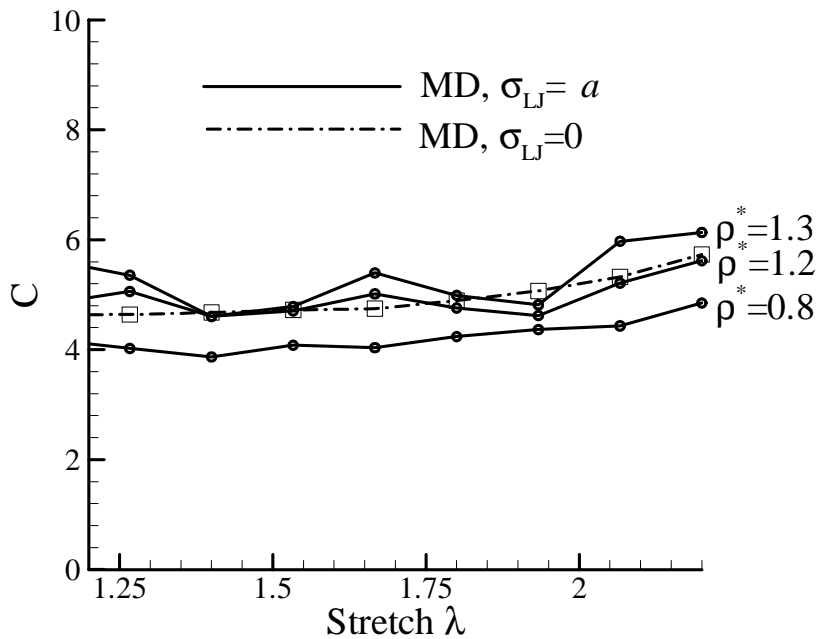
For system that includes EV interactions

$$\langle P_2(\theta_b) \rangle = \frac{F(\rho)}{N_b} (\lambda^2 - \lambda^{-1})$$



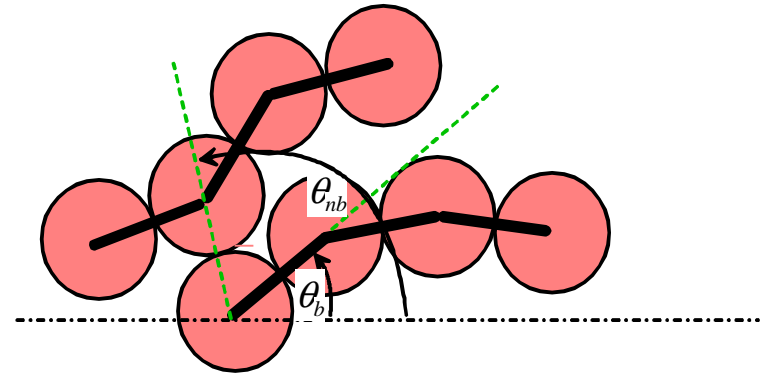
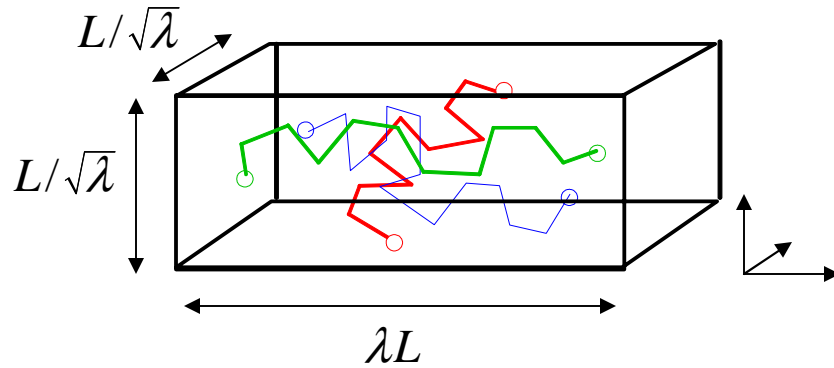
# Relation between stress and bond anisotropy

$$C = \frac{\sigma}{N \langle P_2(\theta_b) \rangle}$$





# Atomic scale contributions to stress



**Virial Theorem** 
$$vt_{ij} = -NkT\delta_{ij} + \sum_{\alpha} \langle f_i^{\alpha} r_j^{\alpha} \rangle$$

**Stress due to covalent bonds** 
$$\sigma_{ij}^b = \frac{1}{kT} \sum_{\alpha \in b} \langle f_i^{\alpha} r_j^{\alpha} \rangle$$

**Stress due to nonbonded interactions**

$$\sigma_{ij}^{nbr} = \frac{1}{kT} \sum_{\alpha \in nbr} \langle f_i^{\alpha} r_j^{\alpha} \rangle$$

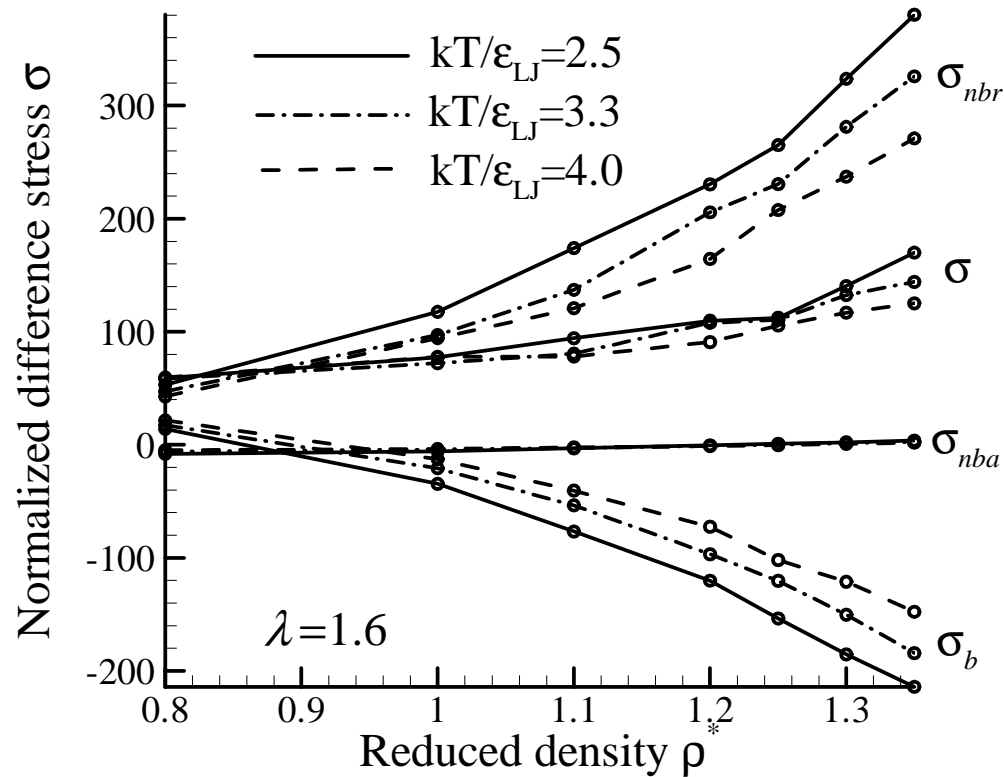
$$\sigma_{ij}^{nba} = \frac{1}{kT} \sum_{\alpha \in nba} \langle f_i^{\alpha} r_j^{\alpha} \rangle$$

$$\sigma^b = \sigma_{11}^b - [\sigma_{22}^b + \sigma_{33}^b]/2$$

$$\Pi^b = \left\{ \sigma_{11}^b(\beta) + \sigma_{22}^b(\beta) + \sigma_{33}^b(\beta) \right\} / 3 \quad \text{etc}$$



# Atomic scale contributions to difference stress

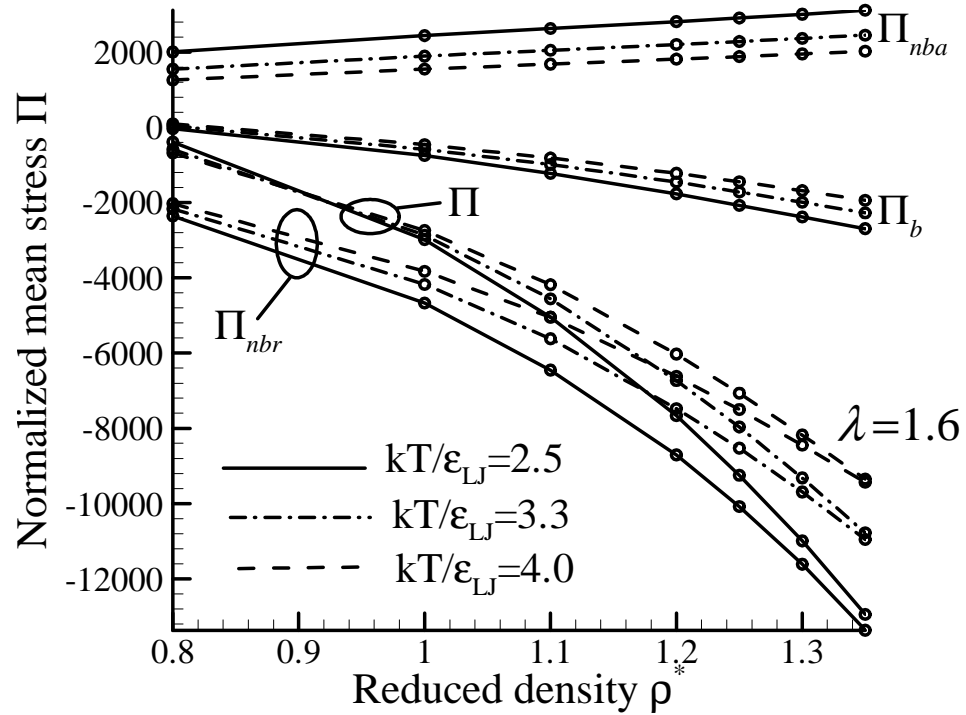


Covalent bonds are in compression, and therefore give a *negative* contribution to difference stress

Dominant contribution to difference stress is from EV interactions

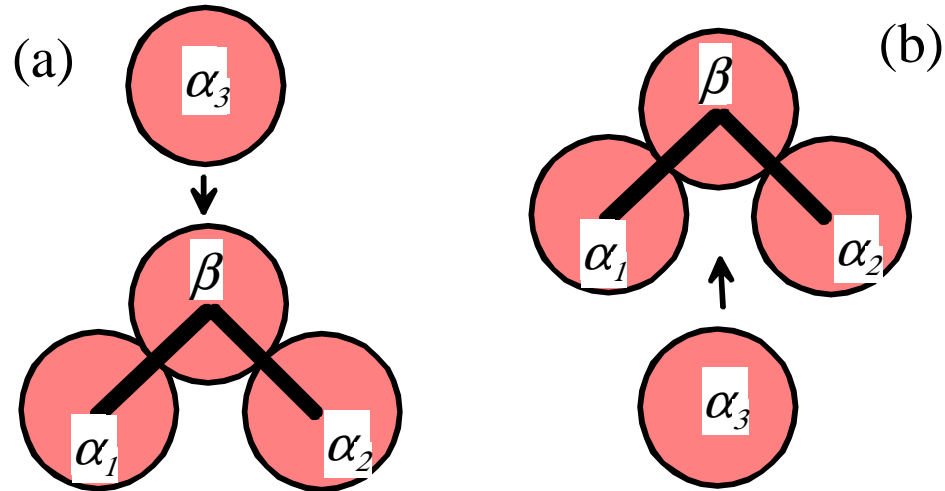


# Atomic scale contributions to mean stress



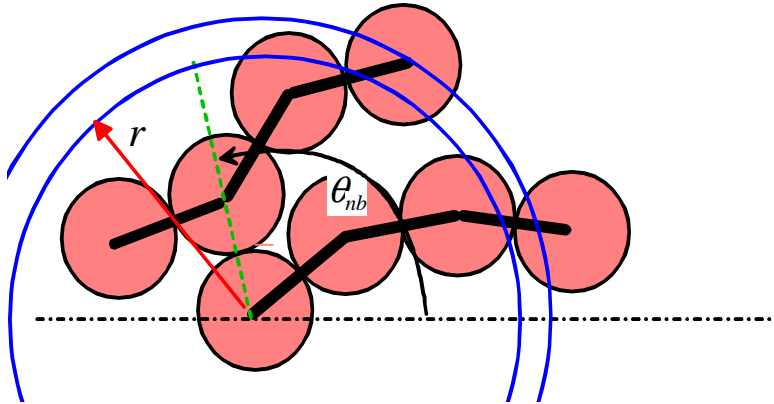


# Steric shielding induces compressive bond forces





# Stress due to repulsive EV interactions



$$dn = \rho g(r) 4\pi r^2 dr$$

$$S_{ij}(r) = \left\langle \frac{r_i r_j}{r^2} \right\rangle$$

$$\langle P_2(\theta_{nbr}) \rangle = \frac{1}{2dn} \sum_{\alpha \in S_r} \left\langle \left\langle (3\cos^2 \theta_\alpha - 1)/2 \right\rangle \right\rangle$$

$$\sigma_{ij}^{nbr} = \frac{1}{kT} \sum_{\alpha \in nb} \langle f_i^\alpha r_j^\alpha \rangle = \frac{2\pi\rho^2 v}{kT} \int_0^\infty u'_{nb}(r) g(r) S_{ij}(r) r^3 dr$$

$$\sigma_{nbr} = \sigma_{11}^{nbr} - [\sigma_{22}^{nbr} + \sigma_{33}^{nbr}] / 2 = \frac{3\pi\rho^2 v}{kT} \int_0^{R_c} u'_{nb}(r) g(r) \langle P_2(\theta_{nbr}) \rangle r^3 dr$$

$$\Pi_{nbr} = [\sigma_{11}^{nbr} + \sigma_{22}^{nbr} + \sigma_{33}^{nbr}] / 3 = \frac{2\pi\rho^2 v}{3kT} \int_0^{R_c} u'_{nb}(r) g(r) r^3 dr$$

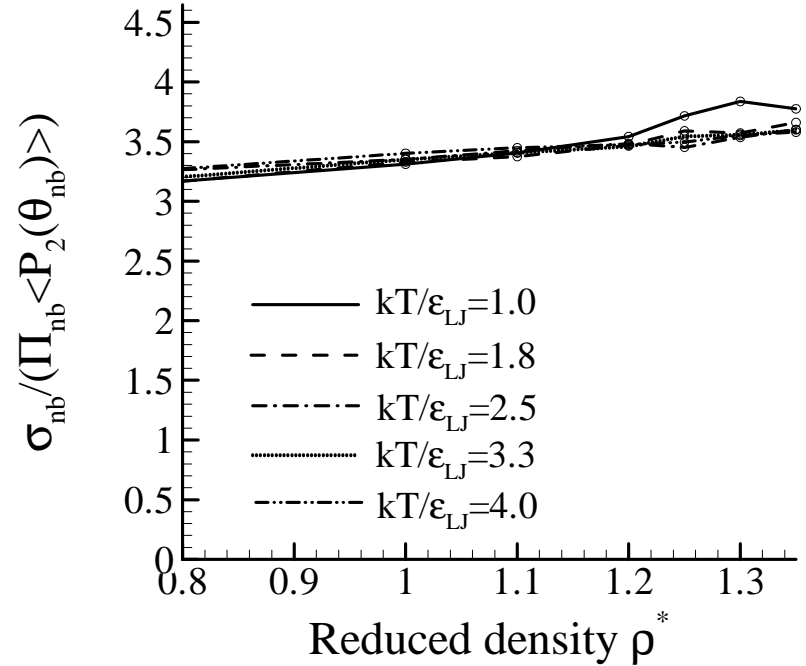
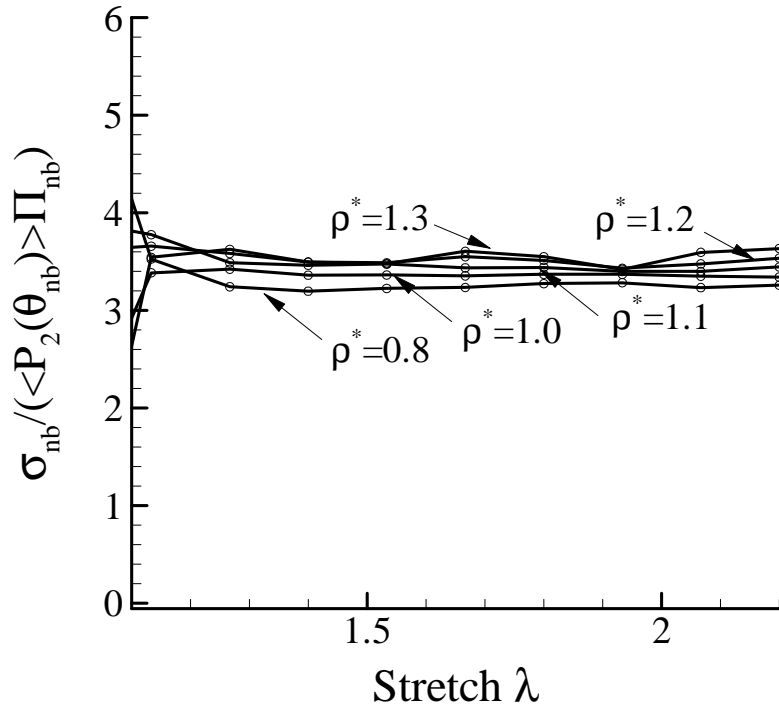
Hard sphere limit

$$\sigma_{nbr} = -3\pi\rho^2 v d^3 g(d) \langle P_2(\theta_{nbr}) \rangle \quad \Pi_{nbr} = -\frac{2}{3}\pi\rho^2 v d^3 g(d)$$

$$\sigma_{nbr} = \frac{9}{2} \Pi_{nbr} \langle P_2(\theta_{nbr}) \rangle$$



# Stress due to repulsive EV interactions

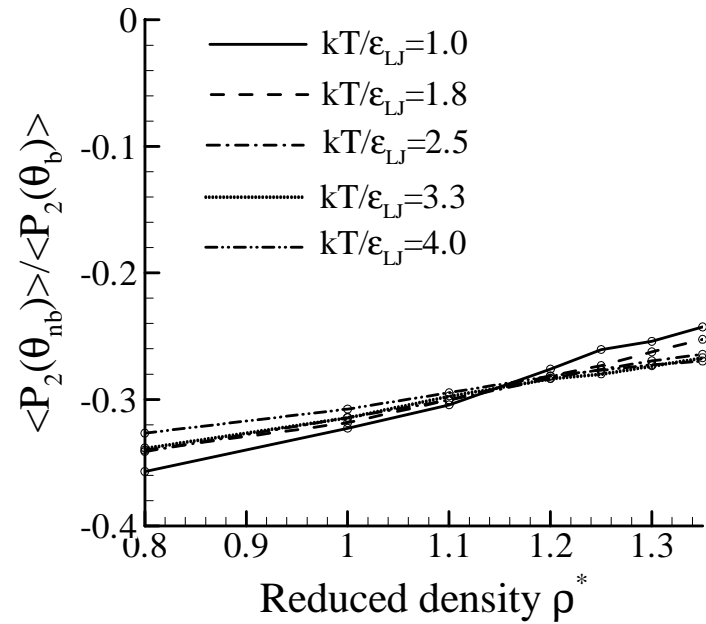
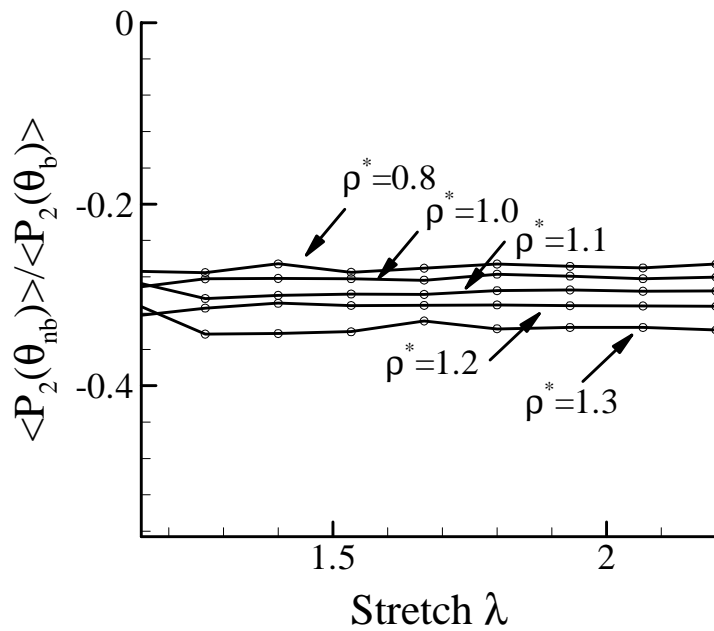
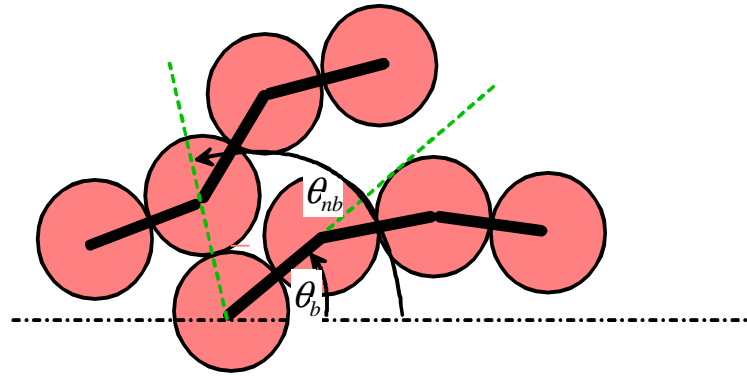


$$\sigma_{nbr} = A \Pi_{nbr} \langle P_2(\theta_{nbr}) \rangle$$

$$A \approx 3.5$$



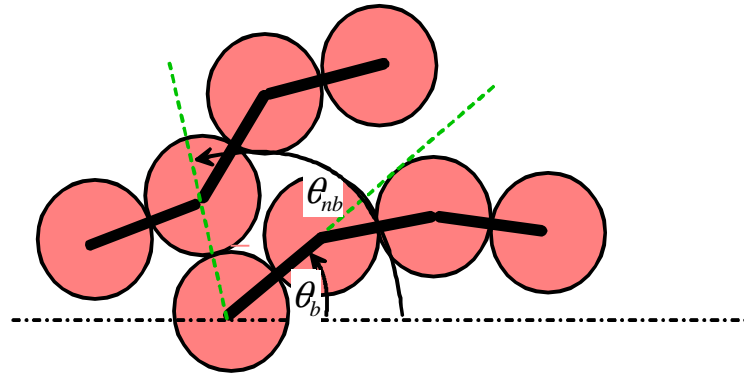
# Orientation of repulsive EV interactions



$$\langle P_2(\theta_{nb}) \rangle = -B \langle P_2(\theta_b) \rangle$$



## Stress due to covalent bonds



$$u_b(r) = \frac{1}{2} \kappa (r - a)^2$$

$$S_{ij}(r) = \left\langle \frac{r_i r_j}{r^2} \right\rangle$$

$$\sigma_{ij}^b = \frac{1}{kT} \sum_{\alpha \in b} \langle f_i^\alpha r_j^\alpha \rangle = \frac{1}{kT} \sum_{\alpha \in b} \left\langle r u_b'(r) \frac{r_i^\alpha r_j^\alpha}{r^2} \right\rangle \approx \frac{1}{kT} \sum_{\alpha \in b} \langle r u_b'(r) \rangle S_{ij}(r)$$

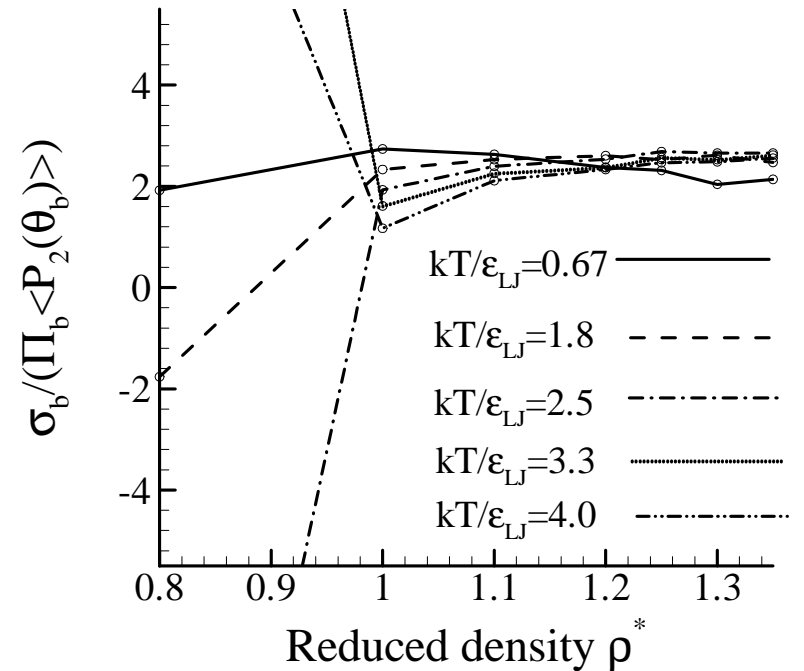
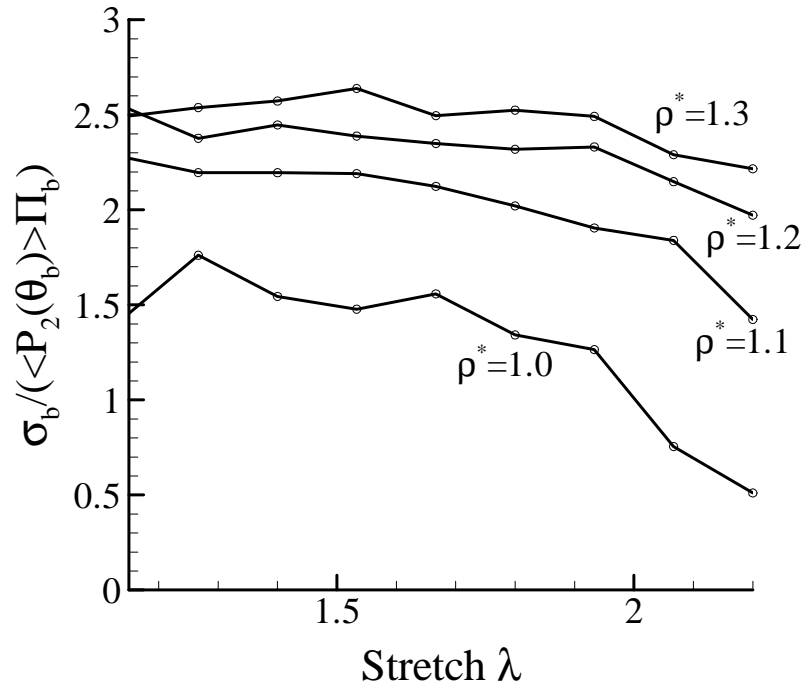
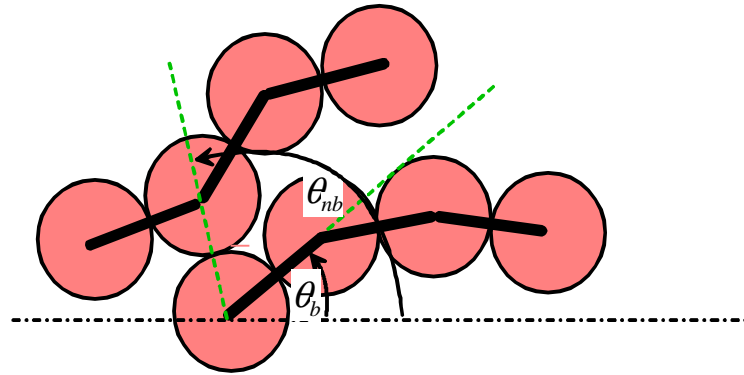
$$\langle r u_b'(r) \rangle = \langle r \kappa (r - a) \rangle = \langle \kappa (r - a)^2 + a \kappa (r - a) \rangle = kT + a \langle f \rangle$$

$$\sigma_b = \sigma_{11}^b - [\sigma_{22}^b + \sigma_{33}^b] / 2 = \left( 1 + \frac{a \langle f \rangle}{kT} \right) \langle P_2(\theta_b) \rangle \quad \Pi_b(\gamma) = \frac{1}{3} \left( 1 + \frac{a \langle f \rangle}{kT} \right)$$

$$\sigma_b = 3 \Pi_b \langle P_2(\theta_b) \rangle$$



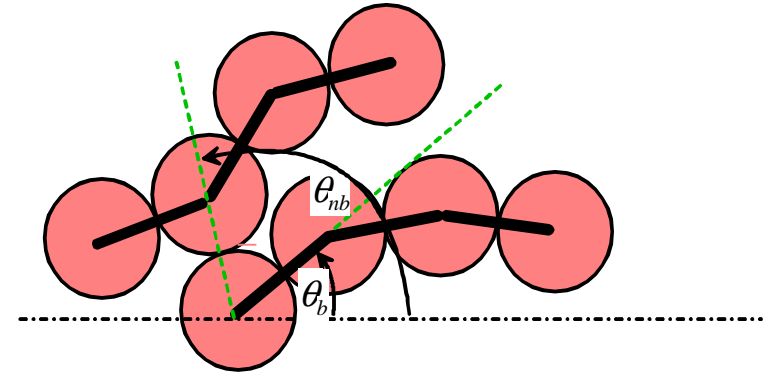
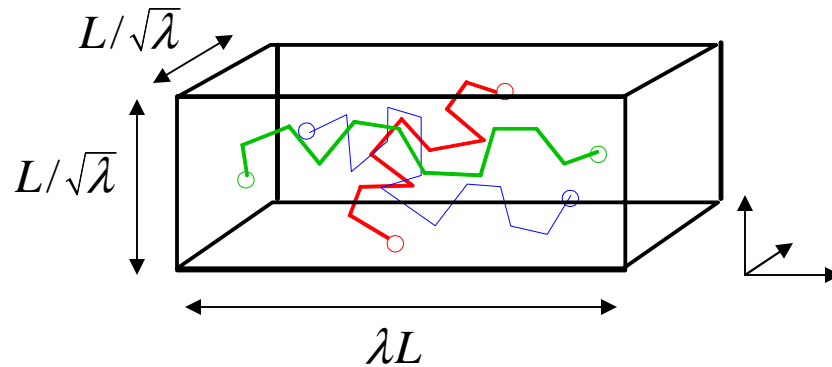
# Stress due to covalent bonds



$$\sigma_b = D \Pi_b \langle P_2(\theta_b) \rangle$$
$$D \approx 2.5$$



# Summary



Covalent bond anisotropy  $\langle P_2(\theta_b) \rangle = \frac{F}{N_b} (\lambda^2 - \lambda^{-1})$

Repulsive EV anisotropy  $\langle P_2(\theta_{nbr}) \rangle = -B \langle P_2(\theta_b) \rangle$

Difference stress  $\sigma = \sigma_{nbr} + \sigma_{nba} + \sigma_b$

Mean stress  $\Pi = -1 + \Pi_{nbr} + \Pi_{nba} + \Pi_b$

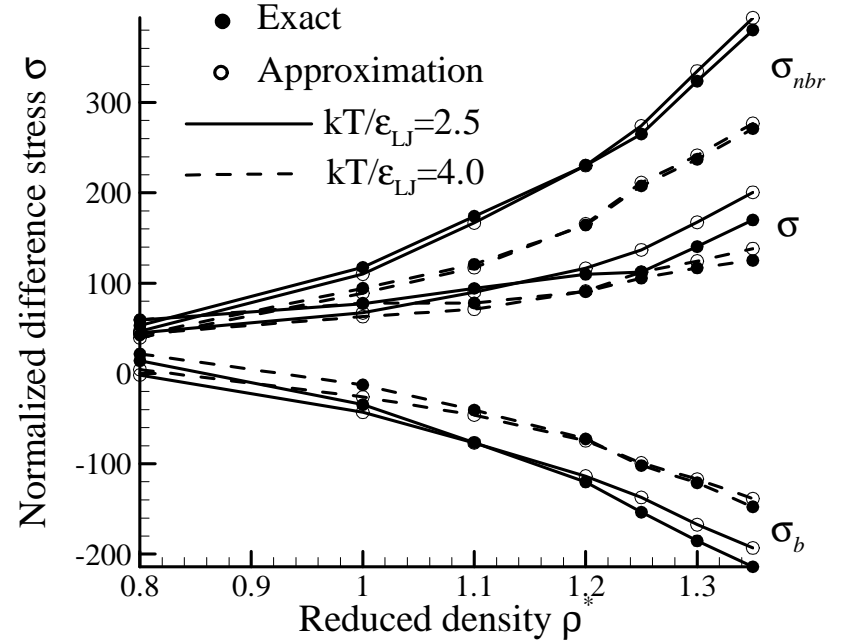
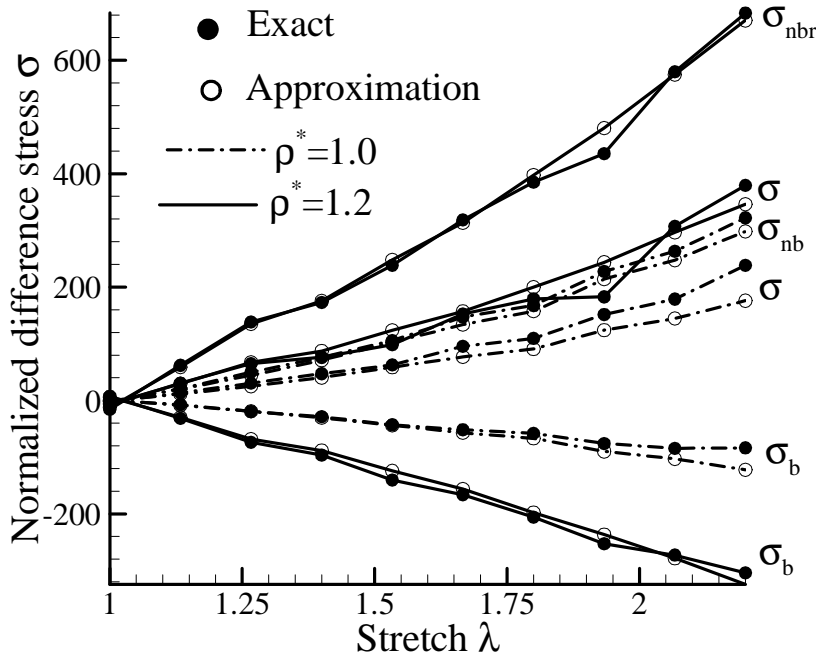
Stress due to covalent bonds  $\sigma_b = D \Pi_b \langle P_2(\theta_b) \rangle$

Stress due to EV interactions  $\sigma_{nba} = 0 \quad \sigma_{nbr} = A \Pi_{nbr} \langle P_2(\theta_{nbr}) \rangle$

**Conclusion**  $\sigma \approx \sigma_b + \sigma_{nbr} = \left( D \Pi_b - 3 A B \Pi_{nbr} \right) \frac{F}{N_b} (\lambda^2 - \lambda^{-1})$



# Comparison of analytical and numerical results



$$\sigma \approx \sigma_b + \sigma_{nbr} = \left( D\Pi_b - AB\Pi_{nbr} \right) \frac{F}{N_b} (\lambda^2 - \lambda^{-1})$$



## Conclusions

- Re-examine atomic scale deformation of elastomers using a model that includes Excluded Volume (EV) interactions
- Deviatoric stress is caused by deformation induced anisotropy of chain network
- Repulsive EV interactions give the dominant contribution to both deviatoric and mean stress
- Pressure and deviatoric stress are intimately connected at atomic scale
- Pressure-deviatoric stress relation  $\sigma \approx \sigma_b + \sigma_{nbr} = \left( D\Pi_b - AB\Pi_{nbr} \right) \frac{F}{N_b} (\lambda^2 - \lambda^{-1})$

## Outstanding issues

- Atomic scale origin of entropic response
- Tension/compression asymmetry
- Role of cross links and junctions
- Glass transition
- Dilation under tensile loading